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Thermal stability of the crystal structure and electronic properties of the high power factor thermoelectric colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$

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ABSTRACT: A complete study on the thermal stability of the crystal structure and electrical properties, as well as on the phase decomposition, of the high performance thermoelectric colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ is presented. Combination of results from differential scanning calorimetry, thermogravimetric analysis, *in situ* neutron powder diffraction, and Seebeck coefficient and electrical resistivity temperature cycling experiments show that the high thermoelectric performance colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ is stable at least up to 700 K in a non-oxidative atmosphere. A superficial sulfur loss without phase decomposition is observed above ≈ 760 K, leading to a progressive increase of both electrical resistivity and Seebeck coefficient. The colusite phase, $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$, starts to decompose above 830 K into Cu_8GeS_6 and CuCrS_2 phases. This leads to a Ge enrichment of the colusite phase (i.e. decrease of the Cu/Ge ratio). This initial decomposition step is followed by a second one above 880 K, related to the structural modification from the cubic structure of colusite into the monoclinic Cu_2GeS_3 structure. Above 930 K, the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ phase is fully decomposed. The influence on the crystal structure of the superficial sulfur loss detected a few dozens of degrees before the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ phase decomposition is discussed and compared to the results obtained on the colusite $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ phase, stable at least up to 1000 K. A singular thermal evolution of the unit cell parameter of the latter is observed and discussed. Finally, a discussion on the thermoelectric performance optimization by process-induced structural disordering in these two colusite phases is given.

INTRODUCTION

In the current global warming crisis, development of alternative technologies to fossil energy such as thermoelectricity, i.e. the conversion of waste heat into electricity, is mandatory. Improvement of the conversion efficiency in line with environmental and cost constraints require the development of new performant thermoelectric materials made of non-toxic and cheap elements. In this context, ternary and quaternary copper-based sulfides, derived from mineral phases, have attracted more and more attention over the past few years: bornite $\text{Cu}_{5-x}\text{T}_x\text{Fe}_{1-y}\text{M}_y\text{S}_4$ ($T = \text{Co}, \text{Zn}; M = \text{Mn}, \text{Zn}; x \leq 0.1; y \leq 0.1$),¹⁻⁷ isocubanite Cu_7FeS_3 ,⁸ famatinite $\text{Cu}_3\text{Sb}_{1-x}\text{M}_x\text{S}_4$ ($M = \text{Ge}, \text{Sn}; x \leq 0.15$),^{9,10} mohite $\text{Cu}_2\text{Sn}_{1-x}\text{M}_x\text{S}_3$ ($M = \text{Co}, \text{Zn}, \text{In}; x \leq 0.25$),¹¹⁻¹³ stannite/kuramite $\text{Cu}_{2+x}\text{Fe}_{1-x}\text{SnS}_4$ ($0 \leq x \leq 1$),^{14,15} kesterite $\text{Cu}_2\text{ZnSnS}_4$,¹⁶ mawsonite $\text{Cu}_6\text{Fe}_2\text{SnS}_8$,¹⁷ stannoidite $\text{Cu}_{8+x}\text{Fe}_{3-x}\text{Sn}_2\text{S}_{12}$ ($x \leq 0.5$),¹⁸ tetrahedrite $\text{Cu}_{12-x}\text{T}_x\text{Sb}_4\text{S}_{13}$ ($T = \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}, \text{Zn}, \text{Mg}; x \leq 2$),¹⁹⁻²⁷ germanite $\text{Cu}_{26-x}\text{Fe}_{4+x}\text{Ge}_4\text{S}_{32}$ ($3.6 \leq x \leq 4.0$),²⁸ renierite $\text{Cu}_{22-x}\text{Zn}_x\text{Fe}_8\text{Ge}_4\text{S}_{32}$ ($1.2 \leq x \leq 2.0$),²⁹ and colusite $\text{Cu}_{26}\text{T}_2\text{M}_6\text{S}_{32}$ ($T = \text{V}, \text{Nb}, \text{Ta}, \text{Cr}, \text{Mo}, \text{W}; M = \text{Ge}, \text{Sn}$).³⁰⁻⁴¹ Among these materials, the synthetic colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ appears to be one of the most promising *p*-type representatives with its high power factor $PF (= S^2/\rho)$, where S is the Seebeck coefficient and ρ is the electrical resistivity) close to $1.5 \text{ mW m}^{-1} \text{ K}^{-2}$ at 300 K and $2.0 \text{ mW m}^{-1} \text{ K}^{-2}$ at 700 K.³⁹ Moreover, these unprecedented electrical

properties among copper-based sulfides are combined with a relatively low thermal conductivity κ , close to $3.0 \text{ W m}^{-1} \text{ K}^{-1}$ at 300 K and $1.6 \text{ W m}^{-1} \text{ K}^{-1}$ at 700 K. This allows this material to reach a figure of merit $ZT (= S^2T/\rho\kappa)$, T being the absolute temperature) close to unity at 700 K and more importantly one of the best average ZT values (close to 0.5 over the 300-700 K temperature range) among this family of materials.³⁹ As described in our previous works,^{39,41} the outstanding electronic transport properties of the synthetic colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ are related to singular interactions between mixed tetrahedral-octahedral complexes $[\text{CrS}_4]\text{Cu}_6$ (arising from the metal-metal interactions between the interstitial Cr atom and the Cu atoms of the host sphalerite framework) and the ionic-covalent conductive “Cu-S” network (i.e. sphalerite framework “ $\text{Cu}_{26}\text{Ge}_6\text{S}_{32}$ ” by omitting Ge atoms). These results revealed that the thermoelectric properties of the synthetic colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ are mostly limited by its thermal conductivity. One possible way to reduce the thermal lattice conductivity is to generate phonon scattering from structural disorder, as demonstrated in synthetic colusite $\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$.³⁶ Indeed, a severe reduction of the lattice thermal conductivity of this material (from $1.4 \text{ W m}^{-1} \text{ K}^{-1}$ down to $0.4 \text{ W m}^{-1} \text{ K}^{-1}$ at 300 K and from $0.6 \text{ W m}^{-1} \text{ K}^{-1}$ down to $0.3 \text{ W m}^{-1} \text{ K}^{-1}$ at 675 K) without affecting its power factor significantly was achieved by controlling the densification process at high temperature and

structural disorder. This leads to an improvement by about 300% of the figure of merit on the full temperature range.³⁶ The reason behind the present thermal stability study is twofold. First, the temperature limit for future potential applications can be determined. Secondly, the feasibility of the temperature-induced disorder approach, successful in colusite $\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$, can be investigated in order to improve the thermoelectric performance of colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$.

Moreover, while some inorganic sulfide compounds are reported with either Mo^{6+} ⁴²⁻⁴⁵ or W^{6+} ⁴⁶⁻⁵² cations in a tetrahedral environment of sulfur, the synthetic colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ is, to the best of our knowledge, the unique representative involving Cr^{6+} cations in such kind of chemical environment. This was assumed to be related to the smaller ionic size and lower electronegativity of Cr^{6+} cation³⁹ that should lead to lower thermal stability of Cr^{6+} containing sulfides compared to those involving Mo^{6+} , W^{6+} or V^{5+} cations. Thus, in addition to the aforementioned applicative interests, the determination of the thermal stability of the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ presents also an interest from the crystal chemistry point of view.

EXPERIMENTAL SECTION

Synthesis. Polycrystalline samples of colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ were synthesized by mechanical alloying and spark plasma sintering following the procedure described elsewhere.³⁹ The resulting pellets were either cut in rectangular ingots to perform temperature cycling measurements of the electrical and thermal properties or ground in agate mortar and sieved down to 150 μm to perform diffraction and DSC/TG analyses.

Differential scanning calorimetry and thermogravimetric analysis. Phase transitions and weight loss were investigated by Differential Scanning Calorimetry (DSC) and ThermoGravimetric Analysis (TGA) respectively, both using a Netzsch STA 449 F3 Jupiter apparatus equipped with a TASC414/4 controller. DSC/TG measurements were carried out on a powdered sample simultaneously from room temperature up to 973 K at the rate of 2 K min^{-1} under argon flow in order to avoid oxidation of the sample. The experimental data were not corrected for the instrumental contributions.

X-ray and neutron powder diffraction. The purity of the samples and the crystal structure of the phases, before and after high temperature characterizations, were checked at room temperature and in air by X-ray powder diffraction (XRPD). High-resolution data were collected using a Bruker D8 Advance Vario 1 two-circle diffractometer (0-2 θ Bragg-Brentano mode) equipped with a $\text{Cu K}\alpha$ X-ray tube, a Ge (111) monochromator (Johansson type) and a LynxEye detector. The structural and chemical stability of the title compound was studied by in situ neutron powder diffraction (NPD) at high temperature at the Institut Laue Langevin (ILL), Grenoble, France. The neutron data were collected from room temperature up to 973 K under dynamical secondary vacuum using the high flux neutron two-axis powder diffractometer D1B ($\lambda = 1.28 \text{ \AA}$) equipped with a one-dimensional curved multidetector. The analysis of the diffraction patterns was performed by Rietveld profile refinement using the FullProf and WinPlotr software packages.^{53,54} NPD refinements were conducted by considering a polynomial function of order 10 for the background description.

Thermoelectric properties measurements. The electrical resistivity (ρ) and Seebeck coefficient (S) were measured simultaneously from ingots using an ULVAC-ZEM3 instrument under partial helium pressure. The samples were cut in the form of parallelepipeds with dimensions of around $3 \times 3 \times 10 \text{ mm}^3$. A first measurement including 7 cycles of heating and cooling (300 K - 700 K - 300 K) was performed, followed by a second measurement

including 7 cycles of heating and cooling (300 K - 760 K - 300 K). Heating and cooling rates were approximately of 2 K min^{-1} .

RESULTS AND DISCUSSION

Sample purity and crystal structure description of colusite. XRPD data acquired on the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ sample after the two-step synthesis process (i.e. mechanical alloying and spark plasma sintering) exhibit the diffraction peaks of the colusite structure only, confirming the absence of secondary phase. Colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ crystallizes in the non-centrosymmetric space group of cubic symmetry $P\bar{4}3n$ with a room temperature unit cell parameter $a = 10.5524(1) \text{ \AA}$. It can be viewed as a superstructure of the cubic sphalerite ZnS structure ($F\bar{4}3m$, $a_s \approx 5.4 \text{ \AA}$) with a $2a_s \times 2a_s \times 2a_s$ supercell induced by an ordered distribution of the metallic cations related to the sphalerite framework “ $\text{Cu}_{26}\text{Ge}_6\text{S}_{32}$ ” on four different crystallographic sites (i.e. 6c, 6d, 8e and 12f) and the occupation of the “interstitial” 2a site (referring to the $P\bar{4}3n$ space group) by Cr^{6+} cations.³⁹ Each cation is in a tetrahedral environment of sulfur anions, the latter being located on two different crystallographic sites (i.e. 8e and 24i). The sphalerite framework “ $\text{Cu}_{26}\text{Ge}_6\text{S}_{32}$ ”, built up of corner-sharing GeS_4 and CuS_4 tetrahedra, hosts the “interstitial” Cr^{6+} cations at the center of a tetrahedron of sulfur sharing edges with the six surrounding CuS_4 tetrahedra (Cu in 12f site). This structural feature, also observed in sylvanite structure and derivative sulfides,^{39,55} leads to abnormally short distances of 2.709 Å between Cr and its six Cu nearest neighbors considered as cations. This implies strong Cr-Cu interactions, which compete with ionic-covalent Cr-S bonds, and consequently to the existence of mixed tetrahedral-octahedral complexes $[\text{CrS}_4]\text{Cu}_6$ isolated from one another.³⁹ It is worth noting that the interactions between these complexes and the conductive “Cu-S” network are responsible for the outstanding electrical transport properties of this colusite.^{39,41}

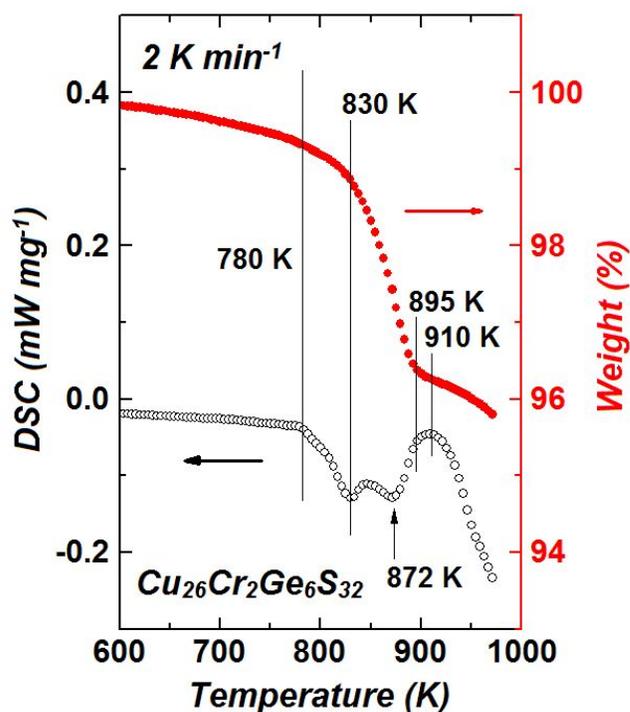


Figure 1. Differential scanning calorimetry and thermogravimetric data for colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ from 600 K up to 973 K at the heating rate of 2 K min^{-1} .

DSC/TGA measurements. Phase transitions and weight loss determined by DSC/TGA measurements from room temperature up to 973 K at the heating rate of 2 K min⁻¹ are shown in Figure 1. Neither signal on the DSC curve nor significant weight loss are detected up to 780 K, indicating that the colusite Cu₂₆Cr₂Ge₆S₃₂ compound is stable up to 780 K in a non-oxidative atmosphere. At higher temperature, both endothermic peaks and weight loss are detected. Thermogravimetric analysis reveals a significant weight loss of about 3 % from 830 K to 895 K, attributed to sulfur volatilization. The DSC curve exhibits a weak endothermic signal starting at 780 K followed by two endothermic peaks with a maximum at 830 K and 872 K respectively and a broad endothermic signal above 910 K. The temperatures for which the sample begins and finishes to lose mass significantly coincide with the maximum of the first endothermic peak (i.e. 830 K) and the end of the second endothermic peak (i.e. 895 K), respectively (Figure 1). However, weight loss continues progressively at higher temperature and reaches a value close to 4 % at around 950 K, corresponding to ≈ 4 atoms of sulfur per formula unit.

High temperature *in situ* neutron powder diffraction measurements. In order to study the structural and chemical stability of the colusite Cu₂₆Cr₂Ge₆S₃₂ compound, *in situ* NPD patterns were recorded every 10 K from 573 K up to 973 K (heating rate of 2 K min⁻¹, 5 minutes per pattern). The neutron wavelength ($\lambda = 1.2847(1)$ Å) was refined using long duration NPD pattern recorded at 293 K by fixing the unit cell parameter of Cu₂₆Cr₂Ge₆S₃₂ to the value refined from room temperature XRPD data ($a = 10.5524(1)$ Å). The refined neutron wavelength was then used for the following neutron data refinements.

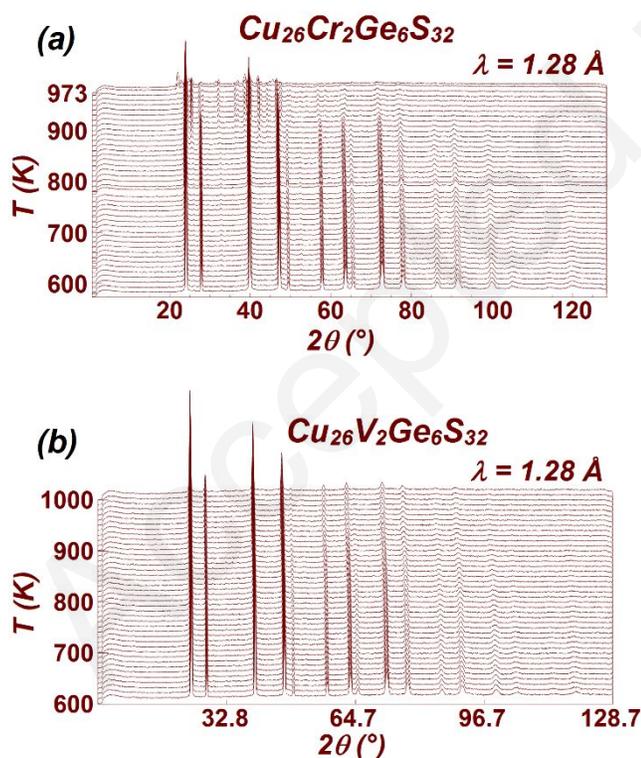


Figure 2. High temperature neutron thermograms of colusites Cu₂₆Cr₂Ge₆S₃₂ (a) and Cu₂₆V₂Ge₆S₃₂ (b) recorded at the heating rate of 2 K min⁻¹.

A qualitative analysis of the high temperature *in situ* NPD patterns suggests a structural stability of the colusite framework up to ≈ 850 K (Figure 2a). Above this temperature, the NPD patterns are characterized by the successive appearance of diffraction peaks arising from secondary phases concomitantly with the vanishing of colusite diffraction peaks. A careful analysis of the NPD patterns reveals that the superstructure diffraction peaks characteristic of the colusite structure are detected up to 880 K, while their existence at higher temperature is not confirmed. Consequently, Rietveld refinements of the NPD data recorded up to 880 K were performed by considering a colusite structure. From these refinements an almost linear thermal evolution of the unit cell parameter of the colusite Cu₂₆Cr₂Ge₆S₃₂ up to 830/840 K could be observed, followed by a relatively constant unit cell parameter and an increase of its standard deviation at higher temperature (Figure 3). This behavior modification is linked to the appearance of diffraction peaks arising from secondary phases at 840 K (Figure 4a) and supported by the huge increase of the χ^2 value above 850 K when only a colusite phase is considered in the refinement (inset of Figure 3). The high quality Rietveld refinement of the NPD pattern recorded at 830 K, shown in Figure 5a, confirms that the sample consists exclusively of the colusite Cu₂₆Cr₂Ge₆S₃₂ phase at this temperature. Moreover, from Rietveld refinements of NPD data recorded up to 830 K, it should be mentioned that (i) the fractional atomic coordinates of the colusite Cu₂₆Cr₂Ge₆S₃₂ phase do not evolve with temperature and (ii) the isotropic displacement parameters of atoms of this phase increase linearly with temperature, with the exception of Cr atom which shows a constant B_{iso} value probably related to its specific chemical environment. Note that similar *in situ* NPD diffraction measurements were performed on a colusite Cu₂₆V₂Ge₆S₃₂ sample. This high temperature characterization indicates the absence of phase decomposition for this colusite up to at least 1000 K (Figure 2b).

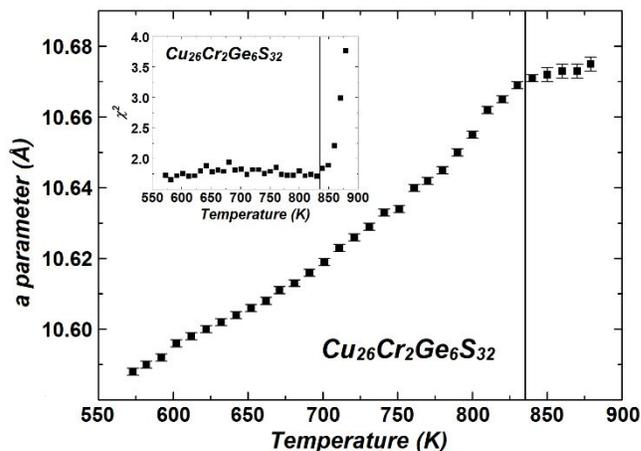


Figure 3. Thermal evolution of the colusite Cu₂₆Cr₂Ge₆S₃₂ unit cell parameter determined from Rietveld refinements of the NPD data. The χ^2 values of the refinements for which only colusite structure was considered are given as an inset.

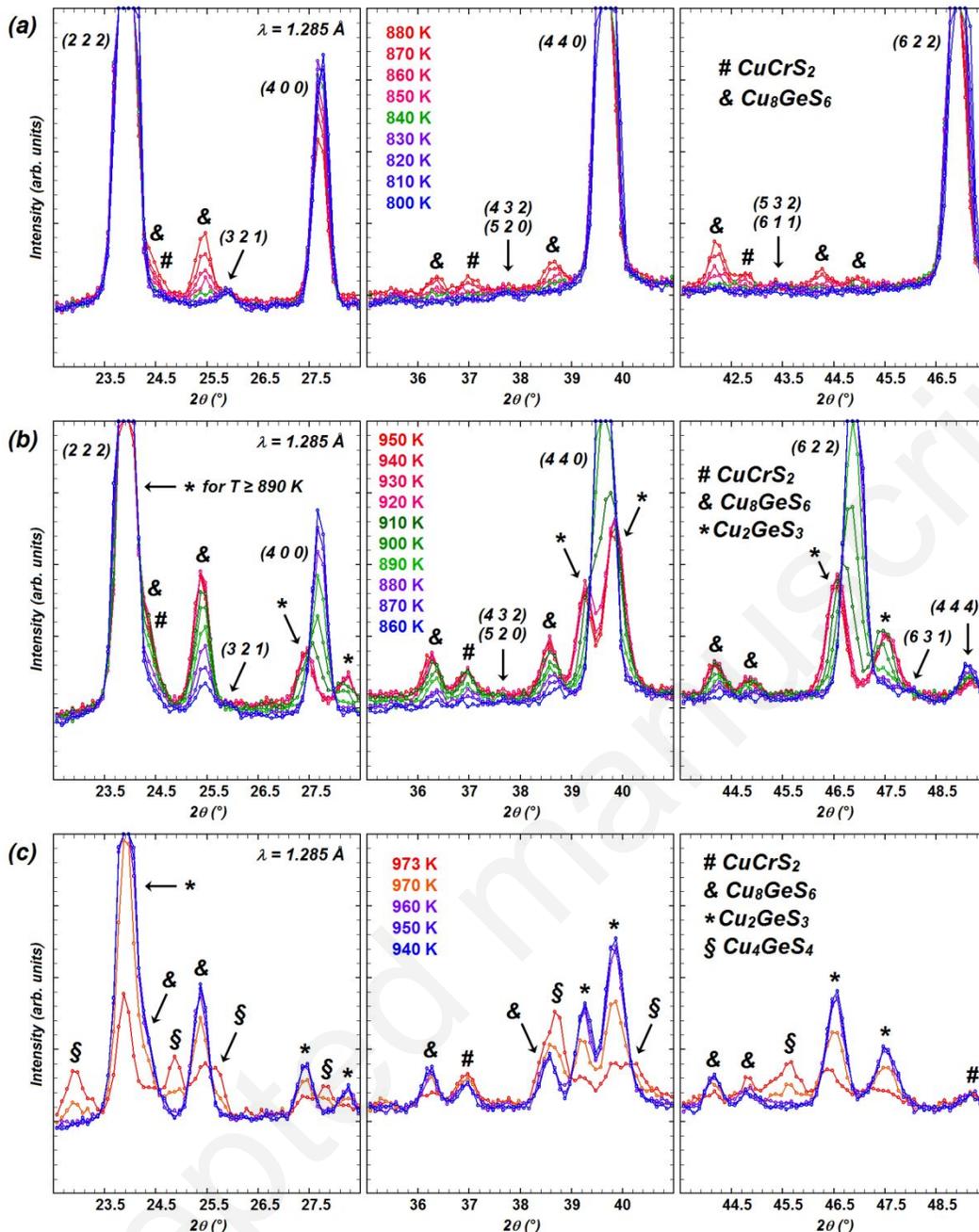


Figure 4. Thermal evolution of the NPD patterns of the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ sample recorded between: (a) 800 K and 880 K; (b) 860 K and 950 K; (c) 940 K and 973 K (heating rate of 2 K min^{-1}). Symbols #, &, *, and § refer to CuCrS_2 phase in its high temperature form, Cu_8GeS_6 in its high temperature cubic form, Cu_2GeS_3 in its monoclinic form, and Cu_4GeS_4 in its monoclinic form, respectively.

As previously mentioned, the in situ NPD study reveals that the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ phase starts to decompose above ≈ 830 K in a non-oxidative atmosphere (Figure 4a). Despite the relatively low intensity diffraction peaks of the phases arising from the thermal decomposition of the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$, it was possible to determine their crystal symmetry and to estimate their chemical composition. Indeed, diffraction peaks observed around 24.5° , 37.0° and 42.6° (Figure 4a) are related to the CuCrS_2 phase, which crystallizes in the trigonal space group $R\bar{3}m$ with unit cell parameters $a = 3.509(2)$ Å and $c = 18.828(19)$ Å at $T = 880$ K (Figure 5b). This high temperature form, reported to be stable above 673 K,⁵⁶ was assumed to be correct and was used in the NPD data refinements. The

remaining diffraction peaks originating from the thermal decomposition of colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ (i.e. around 24.5° , 25.5° , 36.2° , 38.3° , 42.3° , 44.4° , and 44.7° , Figure 4a) can be indexed with the cubic space group $F\bar{4}3m$ and cell parameter $a \approx 10.05$ Å. These results are in agreement with the formation of the high temperature cubic form of the Cu_8GeS_6 phase, a representative of the argyrodite-type chalcogenides family $\text{A}^{\text{m}+}_{(12-n)/\text{m}}\text{B}^{\text{n}+}\text{X}^{2-}_6$, which exhibits a structural phase transition at 328 K from an orthorhombic symmetry to a cubic symmetry.⁵⁷ Consequently, the diffraction peaks arising from the cubic phase of space group $F\bar{4}3m$ have been attributed to the high temperature form of Cu_8GeS_6 . Three structural models have been considered and tested for this phase: Cu_8GeS_6 in

Ag₈GeTe₆-type,⁵⁸ Cu₉GeS_{5,42} in Ag₉GaSe₆-type,⁵⁹ and Cu₈GeS₆ in Ag_{3,3}Cu_{4,7}GeS₆-type⁶⁰ for which Ag atoms have been replaced by Cu atoms. The best refinement was obtained with the latter which corresponds to a putzite mineral derivative structure (Figure 5b). Unfortunately, the exact structural model of the high temperature form of Cu₈GeS₆ cannot be determined from our NPD data.

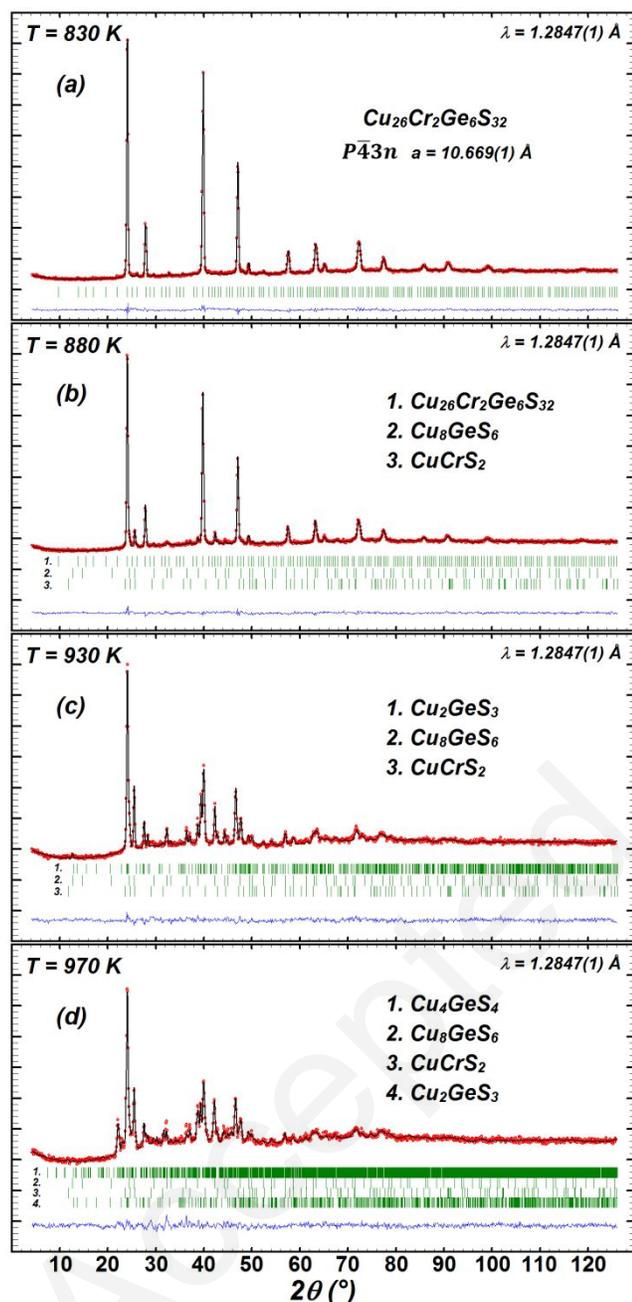


Figure 5. Rietveld refinements of the NPD patterns showing the nature of the phases resulting from the decomposition of the colusite Cu₂₆Cr₂Ge₆S₃₂ at (a) T = 830 K, (b) T = 880 K, (c) T = 930 K, and (d) T = 970 K (heating rate of 2 K min⁻¹).

Hence, this *in situ* NPD investigation suggests that the colusite Cu₂₆Cr₂Ge₆S₃₂ phase is stable up to ≈ 830 K (Figure 5a) and starts to decompose above this temperature into Cu₈GeS₆ and CuCrS₂ phases, both in their high temperature forms. The formation of these two ternary copper sulfides containing germanium or chromium atoms, was confirmed by Rietveld

refinement of the XRPD data recorded at room temperature on the sample after thermal treatment (Figure 6). Both phases were found to crystallize in their room temperature crystal structure: CuCrS₂,⁵⁶ trigonal space group *R3m* with refined unit cell parameters $a = 3.480(1)$ Å and $c = 18.709(1)$ Å; Cu₈GeS₆,⁶¹ orthorhombic space group *Pmn2*₁ with refined unit cell parameters $a = 7.037(1)$ Å, $b = 6.964(1)$ Å and $c = 9.873(1)$ Å.

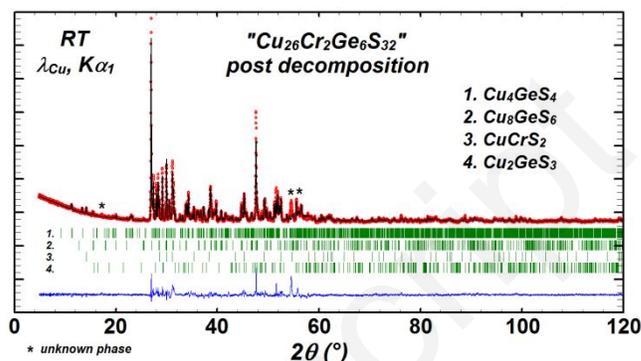


Figure 6. Rietveld refinement of the XRPD pattern recorded at room temperature on sample after thermal treatment showing the nature the phases originating from the decomposition of the colusite Cu₂₆Cr₂Ge₆S₃₂ phase: Cu₄GeS₄ in its monoclinic form (*P2*₁/*c*), Cu₈GeS₆ in its low temperature orthorhombic form (*Pmn2*₁), CuCrS₂ in its low temperature form (*R3m*), Cu₂GeS₃ in its monoclinic form (*Cc*). Symbol * refers to diffraction peaks arising from an unknown phase.

Analysis of the *in situ* NPD data recorded at higher temperature reveals that the intensity of the diffraction peaks of CuCrS₂ and Cu₈GeS₆ phases increases progressively up to ≈ 920 K (Figure 4b), indicating an increase of their weight fractions in the sample. Concomitantly, the diffraction peaks of the colusite Cu₂₆Cr₂Ge₆S₃₂ phase exhibit a rapid decrease in intensity and fully disappear at 920 K. This is evidenced from the thermal evolution of the (4 0 0), (4 4 0) and (6 2 2) reflections (Figure 4b), indicating that, above this temperature, the sample does not contain colusite anymore. Interestingly, from ≈ 890 K up to 920 K, these reflections are splitted into two reflections of relative intensity 2:1 (or 1:2), while the (2 2 2) and (4 4 4) reflections do not evidence any splitting (Figure 4b). Such kind of diffraction peaks splitting suggests modification from a cubic symmetry to a tetragonal symmetry, as already observed in the Cu₂₂₋₃Zn_xFe₈Ge₄S₃₂ series, where a crossover from the cubic germanite-type structure (*P4*₃*n* for $x \leq 0.4$) to the tetragonal renierite-type structure (*P4*₂*c* for $1.2 \leq x \leq 2.0$) was induced by Zn for Cu substitution.²⁹ Indeed, the diffraction peaks of this phase can be indexed with tetragonal unit cell of parameters $a = 10.776$ Å and $c = 10.482$ Å, or $a = 5.388$ Å and $c = 10.484$ Å. However, the diffraction peaks splitting could also be explained by a cubic to monoclinic modification leading to unit cell parameters $a = 6.504$ Å, $b = 11.418$ Å, $c = 6.465$ Å and $\beta = 107.92^\circ$. From these results, several crystal structures have been considered and tested: (i) a renierite derivative structure (space group *P4*₂*c*, $a = 10.776(2)$ Å and $c = 10.482(3)$ Å) for which Fe atoms have been replaced either by Cu or Ge atoms, (ii) a briartite derivative Cu₂FeGeS₄ structure (space group *I4*₂*m*, $a = 5.388(1)$ Å and $c = 10.485(3)$ Å) for which Fe atoms have been replaced by Cu/Ge atoms, and a chalcopyrite derivative CuFeS₂ structure (space group *I4*₂*d*, $a = 5.388(1)$ Å and $c = 10.484(3)$ Å) for which Fe atoms have been replaced by

Ge atoms (on the basis of the work of Ackermann *et al.* on the Ge for Fe substitution of the CuFeS_2 chalcopyrite phase⁶²), and (iii) the crystal structure of the monoclinic Cu_2GeS_3 compound (space group Cc , $a = 6.504(2)$ Å, $b = 11.418(4)$ Å, $c = 6.465(2)$ Å and $\beta = 107.92(2)^\circ$).⁶³ Rietveld refinements performed with the renierite derivatives crystal structure are slightly worse than those done with the other structural models, allowing to exclude the former. Moreover, from the remaining possible crystal structures, that of the monoclinic Cu_2GeS_3 phase leads to a slightly better Rietveld refinement compared to those of the briartite and chalcopyrite derivatives. Consequently, it could be assumed that the phase occurring in the sample above 890 K from the second step of colusite decomposition is the monoclinic Cu_2GeS_3 compound (Figure 5c). It is worth noting that as colusite, its crystal structure is a sphalerite superstructure derivative.

Finally, above ≈ 960 K, the intensity of the diffraction peaks of the Cu_8GeS_6 and Cu_2GeS_3 phases decreases progressively, while the peaks of the CuCrS_2 phase remain almost unchanged, concomitantly with the appearance of new diffraction peaks (labelled with symbol ξ in Figure 4c). The latter are characteristic of the monoclinic Cu_4GeS_4 phase (space group $P2_1/c$), reported to be stable above 605°C (i.e. 878 K).⁶⁴ The formation of this phase is confirmed by Rietveld refinement of the NPD data recorded at 970 K, leading to the refined unit cell parameters $a = 10.051(11)$ Å, $b = 13.365(9)$ Å, $c = 10.033(11)$ Å, and $\beta = 101.36(6)^\circ$ (Figure 5d). The XRPD data recorded at room temperature on the sample after thermal treatment leads to the refined unit cell parameters $a = 9.812(1)$

Å, $b = 13.210(1)$ Å, $c = 9.956(1)$ Å, and $\beta = 101.00(1)^\circ$ (Figure 6). The latter result is in apparent contradiction with the two phase transitions expected for this compound at 339 K and 355 K,⁶⁴ but for which crystal structures have never been reported. Note that in the monoclinic Cu_4GeS_4 phase, Ge, Cu, and S atoms are respectively located on 2, 8, and 8 crystallographic sites of general multiplicity (i.e. $4e$), all fully occupied.⁶⁴ This induces a complex crystal structure characterized by a three-dimensional framework built up of GeS_4 tetrahedra sharing corners with CuS_4 tetrahedra and CuS_3 triangles sharing both corners and edges.⁶⁴ Hence, as with colusites, this sulfide is characterized by the existence of metal-metal interactions.

Temperature cycling influence on the thermoelectric properties. The temperature dependences of the electrical resistivity (ρ) and Seebeck coefficient (S) are displayed in Figure 7 for thermal cycling up to 700 K and 760 K, respectively. In Figures 7a and 7b, seven cycles of heating and cooling performed between 300 K and 700 K confirm the stability of the transport properties, in agreement with DSC and diffraction data. The magnitudes of the electrical resistivity and Seebeck coefficient remain stable during the thermal cycles over the entire temperature range. As exemplified in the inset of Figure 7b, the Seebeck coefficient varies randomly between 86.3 and 87.6 $\mu\text{V K}^{-1}$ at around 340 K. Upon cycling, the power factor is maintained between 1.5 and 1.6 $\text{mW m}^{-1} \text{K}^{-2}$ at 300K and between 1.8 and 1.9 $\text{mW m}^{-1} \text{K}^{-2}$ at 700 K (Figure 7c). Note that these values are identical considering the measurement uncertainty.

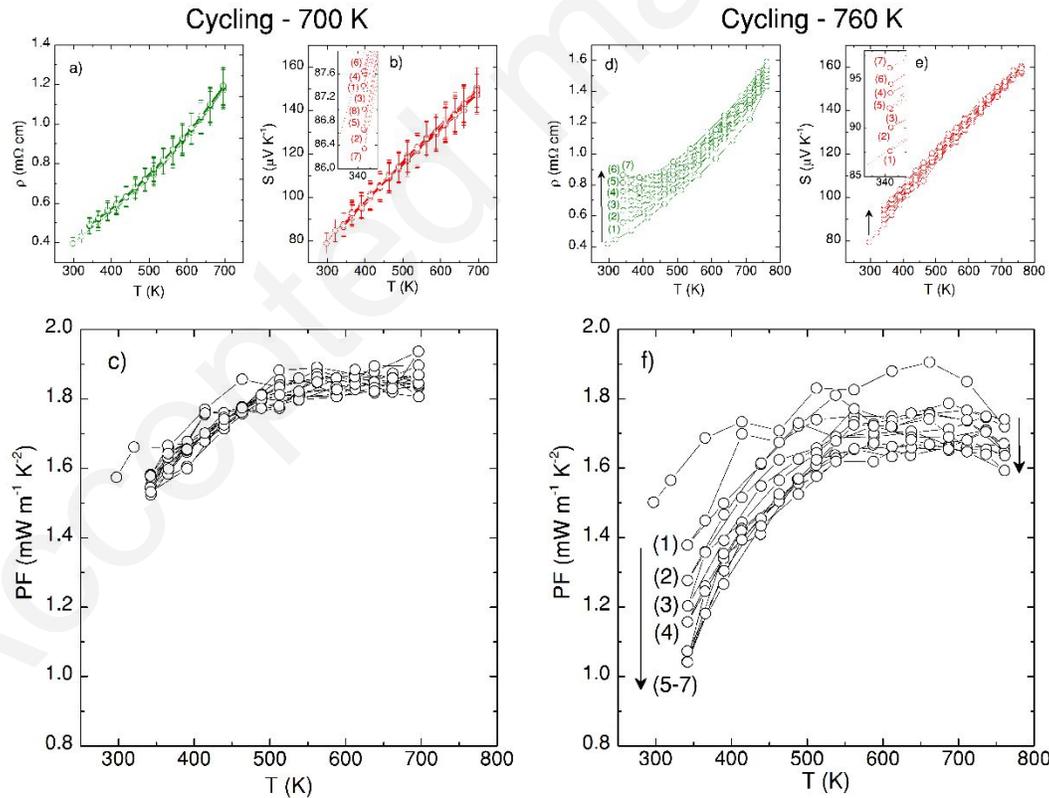


Figure 7. The effect of repeated measurements (7 cycles of heating and cooling) on the electrical resistivity (a, d), Seebeck coefficient (b, e) and power factor (c, f) of $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$; left: from 300 K to 700 K and right: from 300 K to 760 K.

Figures 7d and 7e display the temperature dependences of the electrical resistivity and Seebeck coefficient for thermal cycling between 300 K and 760 K, respectively. With cycling, a slight shift of both the electrical resistivity and Seebeck coefficient towards higher values can be observed. The electrical resistivity at 340 K systematically increases from 0.48 mΩ cm in the first heating ramp to around 0.85 mΩ cm in the 7th cycle. Similarly, the Seebeck coefficient increases from 80 μV K⁻¹ in the 1st cycle to 96 μV K⁻¹ in the 7th cycle. Such behavior likely originates from a decrease in carrier concentration induced by sulfur volatilization above 700 K. Upon cycling, the power factor decreases from 1.6 mW m⁻¹ K⁻² to 1.05 mW m⁻¹ K⁻² at 340 K (Figure 7f).

XRD analysis after cycling at 760 K (Figure 8) confirms that the sample retained high purity (traces of digenite are detected) without significant modification of the colusite unit cell parameter ($a = 10.5524(1)$ Å before cycling and $a = 10.5535(1)$ Å after cycling), suggesting that the sulfur loss remains rather superficial. Interestingly, thermal agitation factors are sensitively smaller than before cycling, which indicates that the repeated cycling process favors the crystallization of the sample. This is easily explained by the relatively short SPS treatment, allowing the subsequent cycling to complete the crystallization.

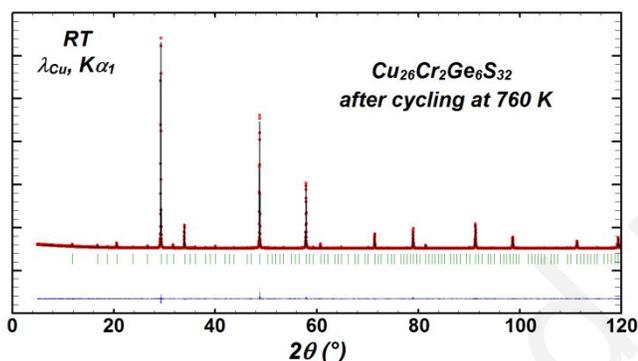
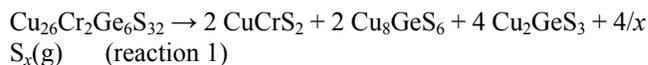
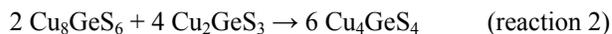


Figure 8. Rietveld refinement of the XRPD pattern recorded at room temperature on sample after cycling at 760 K.

Discussion. *In situ* NPD measurement recorded at the heating rate of 2 K min⁻¹ reveals that the colusite Cu₂₆Cr₂Ge₆S₃₂ phase is stable up to 830 K in a non-oxidative atmosphere (Figure 3 and Figure 4a) and decomposes progressively into Cu₈GeS₆ and CuCrS₂ phases from this temperature. This first step in colusite Cu₂₆Cr₂Ge₆S₃₂ decomposition must induce a germanium enrichment (decrease of the Cu/Ge ratio) of the colusite phase, probably responsible for the second step of the decomposition occurring above 880 K and related to the structural modification from the cubic colusite structure to the monoclinic Cu₂GeS₃ structure (Figure 4b). These temperatures are in fair agreement with those corresponding to the maxima of the endothermic peaks detected at 830 K and 872 K on the DSC curve recorded at the same heating rate (Figure 1). This two-steps phase decomposition is accompanied by a significant weight loss attributed to sulfur volatilization (Figure 1). Moreover, NPD measurement indicates that the colusite phase Cu₂₆Cr₂Ge₆S₃₂ is fully decomposed above 930 K (Figure 4b), with a weight loss of ≈ 4 % (Figure 1), corresponding to ≈ 4 atoms of sulfur per formula unit. Consequently, these results suggest that colusite Cu₂₆Cr₂Ge₆S₃₂ decomposes into CuCrS₂, Cu₈GeS₆ and Cu₂GeS₃ phases, in association with sulfur volatilization, between 830 K and 920 K following the reaction 1:



Finally, above 960 K, the Cu₈GeS₆ and Cu₂GeS₃ phases disappear progressively at the expense of the Cu₄GeS₄ phase (Figure 4c) without significant sulfur loss (Figure 1), following the reaction 2:



Even if the phase decomposition of the colusite Cu₂₆Cr₂Ge₆S₃₂ starts at around 830 K, the thermal evolution of both electrical resistivity and Seebeck coefficient evidenced by cycling measurements performed up to 760 K (Figure 7) suggest a modification of the chemical composition of the colusite phase (i.e. sulfur volatilization) a few dozen degrees before phase decomposition. Note that such an evolution of the electrical properties is not observed during cycling measurements up to 700 K (Figure 7) indicating the high stability of the crystal structure in a non-oxidative atmosphere up to 700 K and consequently of the high performance thermoelectric properties of Cu₂₆Cr₂Ge₆S₃₂ reported in our previous works.^{39,41} A rather superficial sulfur loss at 760 K without decomposition of the colusite Cu₂₆Cr₂Ge₆S₃₂ is supported by the detection of a small endothermic signal on the DSC curve above 780 K as well as weak TGA curve deviation (Figure 1). This superficial sulfur loss likely implies cation over-stoichiometry leading to a metal rich colusite phase with cation occupying initially vacant interstitial crystallographic sites. Moreover, a significant weight loss was evidenced from thermogravimetric measurements in an Ar flow for as-prepared powdered samples of colusites Cu₂₆T₂Sn₆S₃₂ over the temperature range 550-600 K for $T = \text{V}$ and 600-650 K for $T = \text{Nb}, \text{Ta}$.³⁴ This weight reduction was also assigned to volatilization of a part of sulfur implying a larger metal content in sintered samples (1023 K for 1 h under 70 MPa) compared to as-prepared powdered samples. Note that TGA performed on sintered samples do not reveal significant weight loss.

These results show that the colusite Cu₂₆Cr₂Ge₆S₃₂ is highly stable up to 700 K in a non-oxidative atmosphere with a rather superficial sulfur loss around 760 K. The latter leads to crystal structure reorganization and electrical properties evolution without phase decomposition. It is worth noting that these results are obtained from measurements performed under vacuum, which consequently promote sulfur volatilization. On the contrary, sintering processes are performed under an applied pressure of 64 MPa reducing sulfur volatilization and consequently phase decomposition. This explains the high purity pellets prepared by spark plasma sintering even with a sintering temperature of 873 K.³⁹

As with the electrical properties, the superficial sulfur loss between 760 K and 830 K seems to slowly modify the crystal structure of the colusite Cu₂₆Cr₂Ge₆S₃₂. Indeed, no abrupt structural evolution is detected from NPD as evidenced by the almost linear thermal evolution of the refined unit cell parameter of Cu₂₆Cr₂Ge₆S₃₂ from 573 K up to 830 K (Figure 3). From these data, the thermal linear expansion coefficient α_L can be estimated to be $28.9 \times 10^{-6} \text{ K}^{-1}$ for Cu₂₆Cr₂Ge₆S₃₂, a value calculated from the relation $\Delta L/L = \alpha_L \Delta T$ by using a linear fitting of the unit cell parameter between 573 K and 830 K (27 points, $R^2 = 0.9810$) determined from Rietveld refinements of the NPD data. Note that this value is significantly higher than those of $\approx 16.9 \times 10^{-6} \text{ K}^{-1}$ reported for Cu₂₆Nb₂Ge₆S₃₂ between 293 K and 573 K,⁴⁰ and $17.5 \times 10^{-6} \text{ K}^{-1}$ calculated for

$\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$ between 300 K and 650 K, both determined on a lower temperature interval.

As with $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$, we have estimated the thermal linear expansion coefficient α_L of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ at high temperature (between 600 K and 1000 K) from the thermal evolution of its unit cell parameter determined from Rietveld refinements of the NPD data (Figure 9). However, during warming the thermal evolution of the refined unit cell parameter of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ deviates significantly from the linearity hindering the determination of only one thermal linear expansion coefficient. Indeed, an almost linear thermal evolution is observed between 600 K and 830 K ($\alpha_L^{830\text{K}/600\text{K}} = 24.8 \times 10^{-6} \text{ K}^{-1}$, $N = 24$, $R^2 = 0.9856$), followed by an important increase of the a parameter between 830 K and 930 K ($\alpha_L^{930\text{K}/830\text{K}} = 60.1 \times 10^{-6} \text{ K}^{-1}$, $N = 11$, $R^2 = 0.9924$) before slowing down, between 930 K and 1000 K, to a thermal evolution close to that observed between 600 K and 830 K ($\alpha_L^{1000\text{K}/930\text{K}} = 25.7 \times 10^{-6} \text{ K}^{-1}$, $N = 8$, $R^2 = 0.9718$). By considering that the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ is characterized by (i) a superficial sulfur loss between 760 K and 830 K without phase decomposition and (ii) a thermal linear expansion coefficient below 830 K of the same order as that of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ on the same temperature range, it is reasonable to attribute the nonlinear evolution of the thermal expansion of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ observed at high temperature to sulfur volatilization. This is supported by the larger unit cell parameter of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ refined from NPD data recorded at room temperature after the thermal treatment ($a = 10.645(1) \text{ \AA}$) compared to that refined before the thermal treatment ($a = 10.574(1) \text{ \AA}$). Indeed, a similar trend was already observed in $\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$, for which a sample sintered at 1023 K was characterized by a larger unit cell parameter than that refined on a sample sintered at 873 K.³⁶ This difference of cell parameter value was explained by sulfur volatilization promoted by high temperature sintering (i.e. 1023 K instead of 873 K), and subsequently leading to a sample with point defects and disordered regions.³⁶ As with $\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$, the modification of the colusite $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ structure should be expected during warming through sulfur volatilization explaining the nonlinear evolution of its thermal expansion. Nevertheless, artifact from incomplete crystallization of the colusite phase during the fast SPS process cannot be excluded.

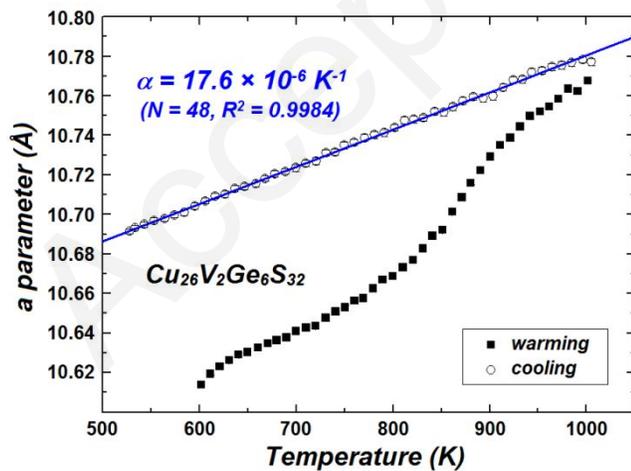


Figure 9. Thermal evolution of the $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ colusite unit cell parameter determined from Rietveld refinements of the NPD data during warming and cooling. Blue line corresponds to the linear fit (see text).

Conversely, the thermal evolution of the a parameter of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ during cooling is perfectly linear over the whole temperature range (Figure 9) and leads to thermal linear expansion coefficient α_L of $17.6 \times 10^{-6} \text{ K}^{-1}$ ($N = 48$, $R^2 = 0.9984$), a value close to those reported for $\text{Cu}_{26}\text{Nb}_2\text{Ge}_6\text{S}_{32}$ and $\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$. The value of the unit cell parameter determined from the extrapolation to 300 K of the linear fit of refined unit cell parameter of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ during cooling ($a = 10.649 \text{ \AA}$) is very close to that determined experimentally ($a = 10.645(1) \text{ \AA}$) from Rietveld refinement of the NPD pattern recorded at room temperature after thermal treatment. This suggests that the crystal structure modification induced during thermal treatment up to 1000 K is stable down to room temperature. Note that potential copper mobility at high temperature, responsible for the nonlinear thermal evolution of the cell parameter in copper-rich tetrahedrite,⁶⁵ is excluded by the irreversibility of the phenomenon during cooling. Consequently, the structural modification induced by superficial sulfur volatilization should lead to different electrical and thermal properties, and subsequently thermoelectric performance, of $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$, as already shown for $\text{Cu}_{26}\text{V}_2\text{Sn}_6\text{S}_{32}$ sintered at different temperatures.³⁶ Such characterizations are in progress.

Finally, the colusite $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ is stable (no phase decomposition) up to at least 1000 K (Figure 2), a temperature well above that of the limit of stability of the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ for which the phase decomposition starts at 840 K (Figure 4a). This indicates that the replacement of pentavalent V^{5+} cations by hexavalent Cr^{6+} cations not only strongly influences the thermoelectric properties³⁹ but also the high temperature stability of the colusite phase. As a consequence, a reduction of the thermal conductivity in order to enhance thermoelectric performance through process controlled structural disordering at high temperature will be strongly limited for $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ and more attainable for $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$. Other strategies must be investigated to induce structural disorder in $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ while maintaining its structural integrity.

CONCLUSION

In this work, we have investigated the effect of temperature on the crystal structure and electrical properties stability of the high thermoelectric performance colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$. The results indicate a high stability of $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ up to at least 700 K in a non-oxidative atmosphere. Above $\approx 760 \text{ K}$ a rather superficial sulfur loss without phase decomposition is evidenced, leading to a crystal structure modification through cation over-stoichiometry. Above 830 K, the colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ phase evidences a two-steps decomposition reaction leading to the formation of Cu_8GeS_6 , CuCrS_2 and Cu_2GeS_3 phases as well as a significant sulfur volatilization of ≈ 4 atoms per formula unit. The colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ phase is fully decomposed above 930 K. Moreover, we show that $\text{Cu}_{26}\text{V}_2\text{Ge}_6\text{S}_{32}$ is stable (no phase decomposition) up to 1000 K, a temperature well above the limit of stability of $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$. This indicates a non-negligible effect of pentavalent V^{5+} by hexavalent Cr^{6+} substitution, not only on the thermoelectric properties, but also on the chemical bonding of this family of material. The direct consequence of this substitution is the difficulty to reduce the thermal conductivity, and subsequently, to enhance the thermoelectric performance of colusite $\text{Cu}_{26}\text{Cr}_2\text{Ge}_6\text{S}_{32}$ through process-controlled structural disordering at high temperature.

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